

EXISTENCE OF LOCAL SUFFICIENTLY SMOOTH SOLUTIONS TO THE COMPLEX MONGE-AMPÈRE EQUATION

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ABSTRACT. We prove the C^∞ local solvability of the n -dimensional complex Monge-Ampère equation $\det(u_{i\bar{j}}) = f(z, u, \nabla u)$, $f \geq 0$, in a neighborhood of any point z_0 where $f(z_0) = 0$.

In the present note we shall consider the complex Monge-Ampère equation

$$(0.1) \quad \det \left(\frac{\partial^2 \phi}{\partial z_i \partial \bar{z}_j} \right) = f(\operatorname{Re} z, \operatorname{Im} z, \phi, \nabla \phi)$$

in an open set Ω of \mathbb{C}^n , f a real-valued function, $z = (z_1, \dots, z_n)$, $z_j = x_j + iy_j$, $\frac{\partial \phi}{\partial z_j} = \frac{1}{2} \left(\frac{\partial \phi}{\partial x_j} - i \frac{\partial \phi}{\partial y_j} \right)$ and $\frac{\partial \phi}{\partial \bar{z}_j} = \frac{1}{2} \left(\frac{\partial \phi}{\partial x_j} + i \frac{\partial \phi}{\partial y_j} \right)$.

Our main result is the following:

Theorem 1. *Let $f \in C^\infty$ be nonnegative near a point $Z^0 = (z_0, u_0, p_0) \in \mathbb{C}^n \times \mathbb{R} \times \mathbb{R}^{2n}$ and $f(Z^0) = 0$. Then for any integer $s > n + 3$, (0.1) has a plurisubharmonic (real-valued) solution $\phi \in H^s$ in a neighborhood of z_0 , such that $\phi(z_0) = u_0$ and $\nabla \phi(z_0) = p_0$ (seen as a function of $(\operatorname{Re} z, \operatorname{Im} z)$).*

Theorem 2. *Suppose in addition to the assumption in Theorem 1 that*

$$\partial_{(x,y)}^\alpha \partial_u^l \partial_p^\beta f(x_0, y_0, u_0, \nabla u_0) = 0$$

for all $|\alpha| + |\beta| + l \leq k - 1$, and there exists $\alpha^* \in \mathbb{N}^{2n}$ such that $|\alpha^*| = k$ and

$$\partial_{(x,y)}^{\alpha^*} f(x_0, y_0, u_0, \nabla u_0) \neq 0.$$

Then (0.1) has a C^∞ plurisubharmonic local real solution u in a neighborhood of (x_0, y_0) .

Recall that in her paper [7] S. Kallel-Jallouli proved the local existence of a smooth plurisubharmonic solution to the problem (0.1) near a point z_0 in the particular case when $f(z, \phi, \nabla \phi) = K(z)g(z, \phi, \nabla \phi)$, $g > 0$, $K(z_0) = 0$ and $dK(z_0) \neq 0$. The proof was essentially based on an inverse function theorem due to G. Nakumara and Y. Maeda [10].

Real Monge-Ampère equations of the form

$$(0.2) \quad \det \left(\frac{\partial^2 \phi}{\partial x_i \partial x_j} \right) = f(x, \phi, \nabla \phi)$$

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were studied by many authors. When f is nonnegative near a point in $\mathbb{R}^n \times \mathbb{R} \times \mathbb{R}^n$, C.S. Lin in [8] proved local existence of a sufficiently smooth solution to (0.2) in the two-dimensional case. The result was then extended in [5] by Hong and Zuily to spaces of arbitrary dimension. The method used is based on a Nash-Moser procedure (see [9]).

In this work, we shall use the same techniques. Since we derive ϕ with respect to two variables in each term of the matrix given in (0.1) and not only just one as in the real case (0.2), some difficulties arise. The introduction of some adapted spaces and the use of some properties of the Laplace operator will enable us to overcome the difficulties.

We will start by giving several fundamental inequalities which play important roles in the proof of convergence of our iteration scheme.

1. LINEAR THEORY

We may assume that $Z^0 = 0$. Following C.S. Lin [8], by means of the change of unknown $\phi = \sum_{i=1}^{n-1} \sigma_i |z_i|^2 + \varepsilon^5 w$ and the change of variables $z_i = \varepsilon^2 \tau_i$, $i = 1, \dots, n$, we can reduce (0.1) to the following equation:

$$(1.1) \quad \det(\phi_{i\bar{j}}) = \det\left((1 - \delta_i^n) \delta_i^j \sigma_i + \varepsilon w_{i\bar{j}}\right) = \tilde{f},$$

where δ_i^j is the Kronecker symbol. The constants σ_i are chosen such that

$$(1.2) \quad \sigma_1 > \sigma_2 > \dots > \sigma_{n-1} = 1.$$

We shall consider

$$(1.3) \quad G(w) = \frac{1}{\varepsilon} \det(\phi_{i\bar{j}}) - \frac{1}{\varepsilon} \tilde{f} \chi(x', y')$$

in the neighborhood of the origin:

$$\Omega = \{(x', x_n, y', y_n) \in \mathbb{R}^{2n}; |x'_i| \leq \pi, |y'_i| \leq \pi, x_n^2 + y_n^2 \leq r^2\},$$

where r will be chosen later and χ is a cut-off function vanishing near $x'_i = \pm\pi$, $y'_i = \pm\pi$ and equal to 1 near the origin.

Note that the projection of Ω on $\mathbb{R}_{(x_n, y_n)}^2$ is the disc

$$D = \{(x_n, y_n) \in \mathbb{R}^2; x_n^2 + y_n^2 \leq r^2\}$$

and the boundary of D is smooth and will help us to use some estimates related to the Laplace operator, established in [4] by Gilbarg and Trudinger.

The linearized operator of G at w is then

$$(1.4) \quad L_G(w) = \sum_{i,j=1}^n \phi^{ij} \partial_i \partial_{\bar{j}} + \sum_{i=1}^n a_i \partial_{\tau_i} + \sum_{i=1}^n b_i \partial_{\bar{\tau}_i} + c,$$

where (ϕ^{ij}) is the matrix of cofactors of $(\phi_{i\bar{j}})$.

Now, for any smooth real-valued function w , the matrix $(\phi_{i\bar{j}})$ is Hermitian and we can find a unitary matrix $T(\tau, \varepsilon)$ satisfying

$$(1.5) \quad T(\tau, \varepsilon) (\phi_{i\bar{j}})^t \bar{T}(\tau, \varepsilon) = \text{diag}(\lambda_1, \dots, \lambda_n).$$

Moreover, we have

Lemma 1.1. *Suppose that w is a smooth real-valued function with $|w|_{C^3} \leq 1$. Then $T(\tau, \varepsilon)$ is smooth in $(\tau, \varepsilon) \in \overline{\Omega} \times [0, \varepsilon_0]$ for some positive ε_0 . Furthermore, with some constant independent of w and ε we have*

$$(1.6) \quad \begin{aligned} & |T_{nn}(\tau, \varepsilon) - 1| + \sum_{i,j=1}^n |\nabla_{x,y} T_{ij}(\tau, \varepsilon)| + \sum_{l=1}^{n-1} |\lambda_l(\tau, \varepsilon) - \sigma_l| \\ & + |\lambda_n(\tau, \varepsilon)| + \sum_{i=1}^{n-1} |T_{in}(\tau, \varepsilon)| \leq c\varepsilon. \end{aligned}$$

Proof. Let $\Phi = (\phi_{i\bar{j}})(\tau, \varepsilon)$. Since $\det(\Phi - \lambda I)(\tau, 0) = -\lambda \sum_{i=1}^{n-1} (\sigma_i - \lambda)$, by (1.2), the eigenvalues $\lambda_i(\tau, \varepsilon)$ are distinct and smooth for small ε and the matrix $T(\tau, \varepsilon)$ is also smooth. It is also clear that $T_{in}(\tau, 0) = 0$ for $i = 1, \dots, n-1$ and $T_{nn}(\tau, 0) = 1$, which implies (1.6). \square

Lemma 1.2. *Let $|w|_{C^2} \leq 1$ and $\theta = \max_{\overline{\Omega}} |G(w)|$. Then the operator*

$$(1.7) \quad -L_G(w) - \theta \Delta,$$

where $\Delta = \sum_{i=1}^n \left(\frac{\partial^2}{\partial x_i^2} + \frac{\partial^2}{\partial y_i^2} \right)$, is a degenerate elliptic operator if ε is small enough.

Proof. We have to prove

$$(1.8) \quad A = \theta |\xi|^2 + \sum_{i,j=1}^n \phi^{ij} \xi_i \bar{\xi}_j \geq 0 \quad \forall (\tau, \xi) \in \Omega \times \mathbb{C}^n.$$

Let us set $\xi = {}^t T(\tau, \varepsilon) \tilde{\xi}$ in (1.8) and let $\tilde{\Phi} = (\phi^{ij})$. We have $A = \theta |\xi|^2 + {}^t \xi \tilde{\Phi} \bar{\xi} = \theta |\xi|^2 + {}^t \tilde{\xi} T \tilde{\Phi} {}^t \bar{T} \tilde{\xi}$ but $\Phi \tilde{\Phi} = \det \Phi \cdot Id$, so by (1.5), $\det \Phi \cdot Id = T \Phi {}^t \bar{T} T \tilde{\Phi} {}^t \bar{T} = \text{diag}(\lambda_i) T \tilde{\Phi} {}^t \bar{T}$, and

$$\begin{aligned} T \tilde{\Phi} {}^t \bar{T} &= \det \Phi \cdot \text{diag} \left(\frac{1}{\lambda_i} \right) = \prod_{i=1}^n \lambda_i \cdot \text{diag} \left(\frac{1}{\lambda_i} \right) \\ &= (\varepsilon G + \chi \tilde{f}) \text{diag} \left(\frac{1}{\lambda_i} \right). \end{aligned}$$

Then,

$$\begin{aligned} A &= \theta \left| \tilde{\xi} \right|^2 + \det \Phi \cdot \sum_{i=1}^n \frac{\left| \tilde{\xi}_i \right|^2}{\lambda_i} \\ &= \theta \left| \tilde{\xi} \right|^2 + \sum_{i=1}^{n-1} \det \Phi \frac{\left| \tilde{\xi}_i \right|^2}{\lambda_i} + \prod_{i=1}^{n-1} \lambda_i \left| \tilde{\xi}_n \right|^2 \\ &= \left(\theta + \prod_{i=1}^{n-1} \lambda_i \right) \left| \tilde{\xi}_n \right|^2 + \sum_{i=1}^{n-1} \frac{\varepsilon G + \tilde{f} \chi + \theta \lambda_i}{\lambda_i} \left| \tilde{\xi}_i \right|^2. \end{aligned}$$

* If $\theta = 0$, then $G(w) = 0$ and $A \geq 0$ if $\tilde{f} \geq 0$.

* If $\theta > 0$, by Lemma 1.1, if ε is small enough, then

$$\varepsilon G + \theta \lambda_i \geq \varepsilon G + \theta \sigma_{n-1} + O(\varepsilon) \geq \frac{1}{2} \theta > 0,$$

and $A \geq 0$.

(1.8) is consequently proved.

Now, we will study a boundary value problem for the degenerate elliptic operator $L_G(w) + \theta \Delta$. First we consider real-valued functions which are periodic in each variable x'_i, y'_i , $1 \leq i \leq n-1$, with period 2π . Following [3], we introduce the space \mathcal{H}_s ($s \in \mathbb{N}$) which is the completion of the space of real-valued trigonometric polynomials $\sum_l \alpha_l(x_n, y_n) e^{i\langle l, (x', y') \rangle}$ with $\bar{\alpha}_l = \alpha_{-l} \in C_0^\infty(\bar{D}(0, r))$ with respect to the norm

$$\|u\|_s^2 = \sum_{t+j \leq s} \sum_l \left(1 + |l|^2\right)^t \|\alpha_l\|_{H^j(\mathring{D}(0, r))}^2.$$

We can define $\mathring{\mathcal{H}}_s$ in the same way, taking α_l in $C_0^\infty(\mathring{D}(0, r))$. For ρ in $\mathring{\mathcal{H}}_1$ we have of course $\rho(x', y', x_n, y_n) = 0$ for $x_n^2 + y_n^2 = r^2$. \square

During this work, we will need two technical lemmas.

Lemma 1.3 (S. L. Sobolev [11]). *If $u \in \mathcal{H}_t$ and $t \geq n + k + 1$, then u is of class C^k , and*

$$\max |\partial^\alpha u| \leq K_{\alpha, t} \|u\|_t, \text{ for } |\alpha| \leq k.$$

$K_{\alpha, t}$ is a constant independent of u .

Lemma 1.4 ([1], [6]). 1) *If $u, v \in L^\infty \cap H^t$ ($t > 0$), then $uv \in L^\infty \cap H^t$ and*

$$\|uv\|_t \leq K_t (\|u\|_{L^\infty} \|v\|_t + \|u\|_t \|v\|_{L^\infty}),$$

where K_t is a constant independent of u, v .

2) *Let $H : \mathbb{R}^m \rightarrow \mathbb{C}$ be a C^∞ function satisfying $H(0) = 0$. If $\omega \in (L^\infty \cap H^s)^m$ ($s > 0$) and $\|\omega\|_{L^\infty} \leq M$, then*

$$\|H(\omega)\|_s \leq K(s, H, M) \|\omega\|_s,$$

where K_s is a constant independent of u

Since each term of the diagonal of the matrix in (0.1) is the Laplace operator applied to $\frac{1}{4}\phi$, we shall need

Lemma 1.5 ([4]). *Let D be a bounded domain in \mathbb{R}^d such that $\partial D \in C^{k+2}$ and $\rho \in H_0^1(D) \cap H^{k+2}(D)$. Then*

$$\|\rho\|_{H^{k+2}(D)} \leq C \left\{ \|\rho\|_{L^2(D)} + \|\Delta \rho\|_{H^k(D)} \right\},$$

where $C = C(d, k, \partial D)$.

The main result of this section is the following.

Theorem 1.6. *Let w be smooth, real-valued, periodic in (x', y') and satisfy the inequality $\|w\|_{C^{n+4}} \leq 1$. Then for any $s_0 \in \mathbb{N}$, one can find a constant $\varepsilon(s_0)$ such that given $g \in C^\infty(\bar{\Omega})$, the problem*

$$(1.9) \quad \begin{cases} L_G(w) \rho + \theta \Delta \rho = g, \\ \rho \in \mathring{\mathcal{H}}_1, \end{cases}$$

admits a unique solution $\rho \in \mathcal{H}_{s_0}$ provided that $0 < \varepsilon < \varepsilon(s_0)$. Moreover, for $0 < s < s_0$, the crucial inequality

$$(1.10_s) \quad \|\rho\|_s \leq C_s \{ \|g\|_s + \|(w)\|_{s+4} \|\rho\|_{L^\infty} \}$$

holds for some constant C_s , independent of w and ε .

Here $\|(w)\|_{s+4}$ is equal to zero if $s \leq n+1$ and to $\|w\|_{s+4}$ if $s > n+1$.

Later, in section 2, we will see how the inequality (1.10) is fundamental and makes the iteration scheme converge to a solution to our problem.

We will divide the proof of Theorem 1.6 into several lemmas. First of all, using the change of unknown function $\underline{\rho} = \rho e^{\lambda|\tau_n|^2}$ with $\tau_n = x_n + iy_n$, we reduce (1.9) to

$$(1.9') \quad \begin{aligned} L(w) \underline{\rho} &= \sum_{i,j=1}^n \left(\phi^{ij} + 4\delta_i^j \theta \right) \partial_i \partial_{\bar{j}} \underline{\rho} + \sum_{i=1}^n a'_i \partial_{\tau_i} \underline{\rho} \\ &+ \sum_{i=1}^n b'_i \partial_{\bar{\tau}_i} \underline{\rho} + c' \underline{\rho} = e^{\lambda|\tau_n|^2} g \end{aligned}$$

with

$$(1.11) \quad \begin{cases} a'_i = a_i - \lambda \tau_n (\phi^{in} + 4\delta_i^n \theta), \\ b'_i = b_i - \lambda \bar{\tau}_n (\phi^{ni} + 4\delta_i^n \theta), \\ c' = c + (\phi^{nn} + 4\theta) (\lambda^2 |\tau_n|^2 - \lambda) - \lambda \bar{\tau}_n a_n - 2\lambda b_n \tau_n. \end{cases}$$

Now, we replace (1.9) by (1.9') and write ρ instead of $\underline{\rho}$. Instead of studying equation (1.9') we will consider the following regularization of (1.9').

Lemma 1.7. *There exist three positive constants $r, \lambda, \varepsilon_0$ such that for $0 < \varepsilon \leq \varepsilon_0$ and any real-valued g in $C^\infty(\bar{\Omega})$ the regularized problem*

$$(1.9'') \quad \begin{cases} L_\nu \rho = L(w) \rho + \nu \Delta \rho = g & \text{in } \Omega, \\ \rho \in \mathring{\mathcal{H}}_1(\Omega), \end{cases}$$

admits a unique (real-valued) solution $\rho \in C^\infty(\bar{\Omega})$ if $\nu > 0$.

Proof. First of all, we recall that as an operator depending on ∂_{x_i} and ∂_{y_i} , $i = 1, \dots, n$, $L_G(w)$ is a real second order operator with real coefficients. Throughout the section $O(\varepsilon)$ means bounded by $K\varepsilon$, where K is a constant independent of ε, λ, r . We shall take $\lambda r = 1$. By (1.3), we can see that $\theta \leq M + O(\varepsilon)$, where M is an absolute constant depending only on the function f . Moreover, by (1.1) we have $\phi^{ij} = O(\varepsilon)$ if $(i, j) \neq (n, n)$, together with its derivatives. We have also $\phi^{nn} = \Upsilon + O(\varepsilon)$ (with $\Upsilon = \prod_{i=1}^{n-1} \sigma_i$), $a_j = O(\varepsilon)$, $b_j = O(\varepsilon)$, $c_j = O(\varepsilon)$. Using (1.11), it follows that

$$\begin{cases} a'_i = O(\varepsilon), & 1 \leq i \leq n-1, \\ b'_i = O(\varepsilon), & 1 \leq i \leq n-1, \\ |a'_n| \leq |a_n| + \lambda(4\theta + \Upsilon)|\tau_n|, & \text{so } a'_n = O(1), \\ |b'_n| \leq |b_n| + \lambda(4\theta + \Upsilon)|\tau_n|, & \text{so } b'_n = O(1), \end{cases}$$

and $c' = (\phi^{nn} + 4\theta) (\lambda^2 |\tau_n|^2 - \lambda) + O(\varepsilon)$, so $-\operatorname{Re} c' \geq \lambda(\Upsilon + 4\theta) (1 - \lambda |\tau_n|^2) + O(\varepsilon)$.

Now, since ρ is 2π -periodic in each variable $x'_i, y'_i, i = 1, \dots, n-1$, and vanishes when $x_n^2 + y_n^2 = r^2$, integration by parts then gives

$$\begin{aligned} (-L_\nu \rho, \rho)_0 &= \int_{\Omega} \nu |\nabla \rho|^2 dx dy \\ &+ \underbrace{\int_{\Omega} \sum_{i,j=1}^n \left(\phi^{ij} + 4\delta_i^j \theta \right) \partial_i \rho \partial_{\bar{j}} \rho dx dy}_{(1)} \\ &- \underbrace{\int_{\Omega} \sum_{i=1}^n a'_i \partial_i \rho \cdot \rho dx dy}_{(2)} - \underbrace{\int_{\Omega} \sum_{i=1}^n b'_i \partial_{\bar{i}} \rho \cdot \rho dx dy}_{(2)'} \\ &- \underbrace{\int_{\Omega} \left(c' + \frac{1}{2} \sum_{i,j=1}^n \partial_{\bar{j}} \partial_i \phi^{ij} \right) \rho^2 dx dy}_{(3)}. \end{aligned}$$

Let us set $\nabla_\tau \rho = {}^t T \tilde{\rho}$, where $\nabla_\tau = \begin{pmatrix} \partial_{\tau_1} \\ \vdots \\ \partial_{\tau_n} \end{pmatrix}$. We get, following the proof of Lemma 1.2,

$$\begin{aligned} (1) &= \int_{\Omega} {}^t (\nabla_\tau \rho) \left(\tilde{\Phi} + 4\theta Id \right) \overline{\nabla_\tau \rho} dx dy \\ &= \int_{\Omega} {}^t \tilde{\rho} T \left(\tilde{\Phi} + 4\theta Id \right) ({}^t \bar{T}) \tilde{\rho} dx dy \\ &= \int_{\Omega} \left\{ \left(4\theta + \prod_{i=1}^{n-1} \lambda_i \right) |\tilde{\rho}_n|^2 + \sum_{i=1}^{n-1} \frac{\varepsilon G + 4\theta \lambda_i + \chi f}{\lambda_i} |\tilde{\rho}_i|^2 dx dy \right\}. \end{aligned}$$

Now, an integration by parts gives

$$\begin{aligned} (2) &= \frac{1}{2} \sum_{i=1}^{n-1} \int_{\Omega} \partial_i a'_i \rho^2 dx dy - \underbrace{\int_{\Omega} a'_n \partial_n \rho \cdot \rho dx dy}_{(4)}, \\ \partial_n \rho &= \sum_{i=1}^n T_{in} \tilde{\rho}_i = T_{nn} \tilde{\rho}_n + \sum_{i=1}^{n-1} \sum_{j=1}^n T_{in} \bar{T}_{ji} \partial_{\tau_j} \rho \\ &= T_{nn} \tilde{\rho}_n + \sum_{i,j=1}^{n-1} T_{in} \bar{T}_{ji} \partial_{\tau_j} \rho + \sum_{i=1}^{n-1} T_{in} \bar{T}_{ni} \partial_{\tau_n} \rho. \end{aligned}$$

Since ${}^t T \bar{T} = Id$, then $\sum_{j=1}^n T_{in} \bar{T}_{ni} = 1$ and

$$(1.12) \quad \partial_n \rho = \frac{1}{|T_{nn}|^2} \left\{ T_{nn} \tilde{\rho}_n + \sum_{i,j=1}^{n-1} T_{in} \bar{T}_{ji} \partial_{\tau_j} \rho \right\}.$$

Now,

$$(4) = - \underbrace{\int_{\Omega} a'_n \rho \frac{T_{nn}}{|T_{nn}|^2} \tilde{\rho}_n dx dy}_{(5)} - \underbrace{\int_{\Omega} a'_n \rho \sum_{i,j=1}^{n-1} \frac{T_{in} \bar{T}_{ji}}{|T_{nn}|^2} \partial_{\tau_j} \rho dx dy}_{(6)}.$$

We have

$$\begin{aligned} (6) &= \frac{1}{2} \int_{\Omega} \sum_{i,j=1}^{n-1} a'_n \frac{T_{in} \bar{T}_{ji}}{|T_{nn}|^2} \partial_{\tau_j} \rho^2 dx dy \\ &= -\frac{1}{2} \sum_{i,j=1}^{n-1} \int_{\Omega} \partial_j \left(a'_n \frac{T_{in} \bar{T}_{ji}}{|T_{nn}|^2} \right) \rho^2 dx dy, \end{aligned}$$

which clearly implies, using Lemma 1.1, that

$$|(6)| \leq C\varepsilon \int_{\Omega} \rho^2 dx dy.$$

For (5) we shall use the inequality $a.b \leq \alpha a^2 + \frac{1}{\alpha} b^2$ with $\alpha = \frac{1}{2} \prod_{i=1}^{n-1} \lambda_i > 0$ and the estimates $a'_n = O(1)$, $T_{nn} = 1 + O(\varepsilon)$. We get

$$\begin{aligned} |(5)| &\leq \int_{\Omega} \frac{|a'_n \rho|}{|T_{nn}|} |\tilde{\rho}_n| dx dy \\ &\leq \frac{1}{2} \int_{\Omega} \prod_{i=1}^{n-1} \lambda_i |\tilde{\rho}_n|^2 dx dy + \int_{\Omega} \frac{2}{\prod_{i=1}^{n-1} \lambda_i} \frac{a_n'^2}{|T_{nn}|} \rho^2 dx dy \\ &\leq \frac{1}{2} \int_{\Omega} \prod_{i=1}^{n-1} \lambda_i |\tilde{\rho}_n|^2 dx dy + C \int_{\Omega} \rho^2 dx dy. \end{aligned}$$

(2)' is estimated similarly to (2). Let us now look at (3). By the discussion at the beginning of the proof we get

$$\operatorname{Re}(3) \geq \left[\lambda(\Upsilon + 4\theta) - (\Upsilon + 4\theta) \lambda^2 |\tau_n|^2 + O(\varepsilon) \right] \int_{\Omega} \rho^2 dx dy.$$

Summing up, we conclude that

$$(-L_{\nu} \rho, \rho) \geq \int_{\Omega} \left\{ \nu |\nabla \rho|^2 + \frac{1}{2} \prod_{i=1}^{n-1} \lambda_i |\tilde{\rho}_n|^2 + [(\Upsilon + 4\theta) \lambda - C + O(\varepsilon)] \rho^2 dx dy \right\} dx dy.$$

We take λ big enough to have coercitivity, and then we apply the Lax-Milgram theorem to get a unique solution ρ in $\mathcal{H}_1(\Omega)$ to the problem (1.9''). Provided g is smooth, the regularity of ρ follows from the theory of elliptic equations. \square

Proof of Theorem 1.6. Suppose (1.10_s) is valid for the regularized problem (1.9'') with a uniform constant C_s for $\nu \in]0, 1]$. Then by letting ν go to zero we shall get a solution of the original problem which of course will satisfy (1.10_s).

For the regularized problem we shall use induction on s . When $s = 0$, (1.10)₀ follows from (1.13). Actually we have the strong estimate

$$\int_{\Omega} |\tilde{\rho}_n|^2 dx dy + \|\rho\|_{L^2}^2 \leq C_0 \|g\|_0^2.$$

Denote by C_s (resp. C) any constant which is independent of ν and ε ; it will change from line to the next. Assume that

$$(1.10'_s) \quad \sum_{\substack{|\alpha| \leq s-1 \\ \alpha_n = \alpha_{2n} = 0}} \int_{\Omega} \left| \left(\widetilde{\partial^\alpha \rho} \right)_n \right|^2 dx dy + \|\rho\|_{s-1}^2 \leq C_{s-1} \left\{ \|g\|_{s-1}^2 + \|(w)\|_{s+3}^2 \|\rho\|_{L^\infty}^2 \right\}$$

for all $\varepsilon \in]0, \varepsilon(s_0)]$ and $0 \leq s-1 \leq s_0$ with $s \geq 1$. We shall find a new $\varepsilon(s_0)$ which makes (1.10'_s) true.

Denote by ρ^s any derivative of order α with $|\alpha| = s$ and $\alpha_n = \alpha_{2n} = 0$. Note that $\partial^\alpha = \partial_{x_1}^{\alpha_1} \dots \partial_{x_{n-1}}^{\alpha_{n-1}} \partial_{y_1}^{\alpha_{n+1}} \dots \partial_{y_{n-1}}^{\alpha_{2n-1}}$

Using (1.13), we get

$$(1.14) \quad |(-L_\nu \rho^s, \rho^s)| \geq C \int_{\Omega} \left[(\rho^s)^2 + \frac{1}{2} \prod_{i=1}^{n-1} \lambda_i \left(\widetilde{\partial^\alpha \rho} \right)_n^2 \right] dx dy.$$

On the other hand,

$$(1.15) \quad (-L_\nu \rho^s, \rho^s) = -(g^s, \rho^s) - ([L_\nu, \partial^\alpha] \rho, \rho^s).$$

Let us look at the second term of (1.15). We have

$$[L_\nu, \partial^\alpha] = \sum_{\substack{\beta \leq \alpha \\ |\beta| \geq 1}} C_{\alpha\beta} [\partial^\beta (\phi^{ij}) \partial_i \partial_j + \partial^\beta a'_j \partial_j + \partial^\beta b'_j \partial_j + \partial^\beta c] \partial^{\alpha-\beta}.$$

For $|\beta| = 1$, we have, using integrations by parts,

$$\begin{aligned} H &= (\partial^\beta (\phi^{ij}) \partial_i \partial_j \rho^{s-1}, \rho^s) = (\partial (\phi^{ij}) \partial_i \partial_j \rho^{s-1}, \partial \rho^{s-1}) \\ &= - \underbrace{\int_{\Omega} \partial_j \rho^{s-1} \partial_i \partial (\phi^{ij}) \partial \rho^{s-1} dx dy}_{(7)} - \int_{\Omega} \partial_j \rho^{s-1} \partial (\phi^{ij}) \partial_i \partial \rho^{s-1} dx dy \\ &= - (7) + \underbrace{\int_{\Omega} \rho^{s-1} \partial_j \partial (\phi^{ij}) \partial_i \partial \rho^{s-1} dx dy}_{(8)} + \int_{\Omega} \rho^{s-1} \partial (\phi^{ij}) \partial \partial_i \partial_j \rho^{s-1} dx dy \\ &= - (7) + (8) - \underbrace{\int_{\Omega} \partial_i \partial_j \rho^{s-1} \partial^2 (\phi^{ij}) \rho^{s-1} dx dy}_{(9)} - H. \end{aligned}$$

Then,

$$\begin{aligned} 2H &= (8) - (7) - (9), \\ (9) &= - \int_{\Omega} \partial_j \rho^{s-1} \partial_i \partial^2 (\phi^{ij}) \rho^{s-1} dx dy - \int_{\Omega} \partial_j \rho^{s-1} \partial^2 (\phi^{ij}) \partial_i \rho^{s-1} dx dy, \end{aligned}$$

which can be estimated by $C\varepsilon \|\rho\|_s^2$. We estimate the terms (7), (8) in the same way to obtain

$$|H| \leq C\varepsilon \|\rho\|_s^2.$$

For $2 \leq |\beta| \leq n+1$, since $\|w\|_{C^{n+3}} \leq 1$ and $|\alpha| - |\beta| + 2 \leq s$, we have

$$(1.16) \quad \left| \left(\partial^\beta (\phi^{ij}) \partial_i \partial_j \rho^{s-|\beta|}, \rho^s \right) \right| \leq C\varepsilon \|\rho\|_s^2.$$

If $|\beta| > n + 1$, we can write

$$A = \left| \left(\partial^\beta (\phi^{ij}) \partial_i \partial_{\bar{j}} \rho^{s-|\beta|}, \rho^s \right) \right| \leq \|\rho\|_s \left\| \partial^\beta (\phi^{ij}) \partial_i \partial_{\bar{j}} \rho^{s-|\beta|} \right\|_{L^2}.$$

Using Lemma 1.4 with $u = \partial^2 (\phi^{ij})$ and $v = \rho$, we get

$$(1.17) \quad A \leq C \|\rho^s\|_{L^2} (\varepsilon \|w\|_{C^4} \|\rho\|_s + \|\rho\|_{L^\infty} \|w\|_{s+4}).$$

The lower order terms in L_ν are estimated in a similar way. Summing up, we conclude, using (1.15) to (1.17), that

$$(1.18) \quad |(-L_\nu \rho^s, \rho^s)| \leq C \left\{ \varepsilon \|\rho\|_s^2 + \|\rho^s\|_{L^2} [\|g^s\|_{L^2} + \|(w)\|_{s+4} \|\rho\|_{L^\infty}] \right\}.$$

It remains to estimate the derivatives $\partial_{x_n}^{k_1} \partial_{y_n}^{k_2} \rho^{s-k_1-k_2}$, for $k_1 + k_2 = 1, \dots, s$. From (1.12) we get

$$\partial_n \rho^{s-1} = \frac{1}{|T_{nn}|^2} \left\{ T_{nn} (\widetilde{\rho^{s-1}})_n + \sum_{i,j=1}^{n-1} T_{in} \bar{T}_{ji} \partial_j \rho^{s-1} \right\},$$

but since ρ is real-valued, then

$$\partial_{x_n} \rho^{s-1} = \operatorname{Re} (\partial_{\tau_n} \rho^{s-1}) \quad \text{and} \quad \partial_{y_n} \rho^{s-1} = -\operatorname{Im} (\partial_{\tau_n} \rho^{s-1}),$$

so

$$\begin{cases} \|\partial_{x_n} \rho^{s-1}\|_{L^2} \leq C \left(\left\| (\widetilde{\rho^{s-1}})_n \right\|_{L^2} + \varepsilon \|\rho\|_s \right), \\ \|\partial_{y_n} \rho^{s-1}\|_{L^2} \leq C \left(\left\| (\widetilde{\rho^{s-1}})_n \right\|_{L^2} + \varepsilon \|\rho\|_s \right). \end{cases}$$

Using the induction estimate (1.10'_{s-1}), we get, for $k_1 + k_2 = 1$,

$$(1.19) \quad \|\partial_{x_n}^{k_1} \partial_{y_n}^{k_2} \rho^{s-1}\|_{L^2} \leq C \{ \|g\|_{s-1} + \|(w)\|_{s+3} \|\rho\|_{L^\infty} + \varepsilon \|\rho\|_s \}.$$

Denote $\partial_n^k = \partial_{x_n}^{k_1} \partial_{y_n}^{k_2}$ for $k = k_1 + k_2$ and write the original équation (1.9'') in another way:

$$(1.20) \quad \begin{aligned} \Delta_n \rho &= \frac{4}{\phi^{nn} + 4\theta + 4\nu} \left\{ g - \sum_{i \leq n} \left(\frac{1}{4} \phi^{ii} + \theta + \nu \right) \Delta_i \rho - \sum_{i \neq j} \phi^{ij} \partial_i \partial_{\bar{j}} \right. \\ &\quad \left. - a'_j \partial_{\tau_j} \rho - b'_j \partial_{\bar{\tau}_j} \rho - c' \rho \right\} = F, \end{aligned}$$

with $\Delta_i = \partial^2 / \partial x_i^2 + \partial^2 / \partial y_i^2$, $i = 1, \dots, n$.

Applying ∂^{s-1} to (1.20), we obtain using Lemma 1.5, with $d = 2$ and $D = \{(x_n, y_n); x_n^2 + y_n^2 \leq r^2\}$,

$$(1.21) \quad \|\partial_n^{k+2} \rho^{s-k}\|_{L^2}^2 \leq C \left\{ \|\rho^{s-k}\|_{L^2}^2 + \sum_{0 \leq l \leq k} \|\partial_n^l \partial^{s-k} F\|_{L^2}^2 \right\}.$$

Now, we are able to prove (1.19) for $k = 2, \dots, s$ by induction on k , using (1.21), (1.10'_{s-1}) and Lemma 1.4 (1.25). We get

$$(1.22) \quad \sum_{k=2}^s \|\partial_n^k \rho^{s-1}\|_{L^2} \leq C_s \{ \|g\|_s + \|(w)\|_{s+3} \|\rho\|_{L^\infty} + \varepsilon \|\rho\|_s \}.$$

Using (1.14), (1.18), (1.19) and (1.22), we get (1.10'_s) for $0 \leq s \leq s_0$ and $0 \leq \varepsilon \leq \varepsilon(s_0)$ if $\varepsilon(s_0)$ is small enough. \square

2. EXISTENCE OF SUFFICIENTLY SMOOTH SOLUTIONS

In this section we shall construct, using the results of section 1, a sequence which converges to a solution to our problem.

Let E be the space of all smooth functions in $\overline{\Omega}$ which are periodic in (x', y') , and let τ and $\sigma > 1$ be two constants which will be chosen later. One can define a family of smoothing operators S_n from E to E such that with $\mu_n = \sigma^{\tau^n}$ and for $0 \leq s_1, s_2 \leq s_0$

$$(2.1) \quad \|S_n u\|_{s_1} \leq C_{s_2} \|u\|_{s_2}, \quad \text{if } s_1 \leq s_2,$$

$$(2.2) \quad \|S_n u\|_{s_1} \leq C_{s_1} \mu_n^{s_1-s_2} \|u\|_{s_2}, \quad \text{if } s_1 \geq s_2,$$

$$(2.3) \quad \|S_n u - u\|_{s_1} \leq C_{s_2} \mu_n^{s_1-s_2} \|u\|_{s_2}, \quad \text{if } s_1 \leq s_2,$$

where the constant C_{s_i} is independent of u , n , σ , τ . We will construct w_n , $n = 0, 1, \dots$, by induction on n as follows. Starting with $w_0 = 0$, suppose w_0, w_1, \dots, w_n have been chosen and define w_{n+1} as follows:

$$(2.4) \quad w_{n+1} = w_n + S_n \rho_n,$$

where ρ_n is the solution of

$$(2.5) \quad L(w_n) \rho_n = L_G(w_n) \rho_n + \theta_n \triangle \rho_n = g_n \text{ in } \Omega$$

satisfying

$$(2.6) \quad \rho_n \in \overset{\circ}{\mathcal{H}}_1$$

given by Theorem 1.6, when

$$(2.7) \quad g_n = -G(w_n), \quad \theta_n = \sup_{\overline{\Omega}} |G(w_n)|.$$

In order to ensure that the w_k are well defined, there are several things to be verified. We prove first the following result.

Lemma 2.1. *Suppose $\|w_k\|_{C^{n+3}} \leq 1$ for $k = 0, 1, \dots, m$. Then for all $k = 0, 1, \dots, m$ we have*

$$(2.8) \quad \|g_k\|_s \leq C_s \{ \|g_0\|_s + \|w_k\|_{s+2} \},$$

$$(2.9) \quad \|w_{k+1}\|_{s+4} \leq C_s^{k+1} \mu_{k+1}^\beta \|g_0\|_s \quad \text{for some } \beta > \frac{4}{\tau-1},$$

$$(2.10) \quad \|g_{k+1}\|_{L^2} \leq \mu_{k+1}^{-\kappa} \|g_0\|_{s^*} \quad \text{for some positive constants } \kappa \text{ and } s^*.$$

Proof. a) $\|g_k\|_s \leq \|G(w_0)\|_s + \|G(w_k) - G(w_0)\|_s$. Using the Taylor formula for G , Lemma 1.4 and the hypothesis $w_0 = 0$, we easily get (2.8).

b) Using (2.2) and (2.4), we get

$$\|w_{k+1}\|_{s+4} \leq C_s (\|w_k\|_{s+4} + \|S_k \rho_k\|_{s+4}) \leq C_s (\|w_k\|_{s+4} + \mu_k^4 \|\rho_k\|_s).$$

* If $s \leq n+1$, (1.10_s) gives

$$\|\rho_k\|_s \leq C_s \|g_k\|_s \leq C' \{ \|g_0\|_s + \|w_k\|_{s+2} \},$$

so, since $\mu_k > 1$, then

$$(2.11) \quad \|w_{k+1}\|_{s+4} \leq C_s \mu_k^4 (\|g_0\|_s + \|w_k\|_{s+4}).$$

* If $s > n+1$, we have by Lemma 1.3,

$$\|\rho_k\|_{L^\infty} \leq C \|\rho_k\|_{n+1},$$

so (1.10_s) gives $\|\rho_k\|_s \leq C(\|g_k\|_s + \|w_k\|_{s+4})$, and using (2.8) we get the same estimate (2.11). Iteration of this estimate yields, since $w_0 = 0$,

$$\|w_{k+1}\|_{s+4} \leq C_s^k (k+1) \mu_0^4 \dots \mu_k^4 \|g_0\|_s \leq C_s^k (k+1) \sigma^4 \frac{\tau+1}{\tau-1} \|g_0\|_s,$$

which proves (2.9).

c) In view of (2.5) and (2.7) we have

$$-g_{k+1} = G(w_k) + L_G(w_k) S_k \rho_k + Q(w_k, S_k \rho_k),$$

where Q is the quadratic error, and

$$-g_{k+1} = G(w_k) + L(w_k) \rho_k + L(w_k) (S_k - I) \rho_k - \theta_k \triangle S_k \rho_k + Q(w_k, S_k \rho_k).$$

The sum of the first two terms in the right-hand side vanishes, so

$$\|g_{k+1}\|_0 \leq \underbrace{\|L(w_k) (S_k - I) \rho_k\|_0}_{(1)} + \underbrace{\|\theta_k \triangle S_k \rho_k\|_0}_{(2)} + \underbrace{\|Q(w_k, S_k \rho_k)\|_0}_{(3)}.$$

Since $\|w_k\|_{n+1} \leq 1$, then

$$\begin{aligned} |(1)| &\leq C \|(S_k - I) \rho_k\|_2 \leq C_{s^*} \mu_k^{-(s^*-2)} \|\rho_k\|_{s^*}, \\ |(2)| &\leq \theta_k \|S_k \rho_k\|_2 \leq \theta_k \mu_k^2 \|\rho_k\|_0, \\ |(3)| &\leq C \|S_k \rho_k\|_{L^\infty} \|S_k \rho_k\|_4. \end{aligned}$$

It follows from (1.10_{s*}), with s^* to be determined, that

$$\|g_{k+1}\|_0 \leq C_{s^*} \left\{ \mu_k^{-(s^*-2)} (\|g_k\|_{s^*} + \|w_k\|_{s^*+4}) + \theta_k \mu_k^2 \|g_k\|_0 + \mu_k^{n+5} \|g_k\|_0^2 \right\}.$$

Now by (2.8) and (2.9)

$$\|g_k\|_{s^*} + \|w_k\|_{s^*+4} \leq C(\|g_0\|_{s^*} + \|w_k\|_{s^*+4}) \leq C C_{s^*}^k \mu_k^\beta \|g_0\|_{s^*},$$

so

$$(2.12) \quad \|g_{k+1}\|_0 \leq C_{s^*} \left\{ \mu_k^{-(s^*-2-\beta)} C^k \|g_0\|_{s^*} + \theta_k \mu_k^2 \|g_k\|_0 + \mu_k^{n+5} \|g_k\|_0^2 \right\}.$$

Now by (2.7), and using Lemma 1.3, we get

$$\theta_{k+1} = \|g_{k+1}\|_{L^\infty} \leq \|g_{k+1}\|_{n+1},$$

so

$$\theta_{k+1} \leq \|L(w_k) (S_k - I) \rho_k\|_{n+1} + \theta_k \|\triangle S_k \rho_k\|_{n+1} + \|Q(w_k, S_k \rho_k)\|_{n+1}.$$

Using the same estimates as before, we get

$$(2.13) \quad \theta_{k+1} \leq C_{s^*} \left\{ \mu_k^{-(s^*-n-3-\beta)} C_{s^*}^k \|g_0\|_{s^*} + \theta_k \mu_k^{n+3} \|g_k\|_0 + \mu_k^{6+2n} \|g_k\|_0^2 \right\}.$$

Let us choose κ and s^* such that

$$(2.14) \quad \begin{cases} (2-\tau)\kappa - 8 - 2n > 0, \\ s^* - 4 - n - \beta - \kappa\tau > 0, \end{cases}$$

and set

$$(2.15) \quad d_{k+1} = \max(\mu_{k+1}^\kappa \|g_{k+1}\|_0, \mu_{k+1}^\kappa \theta_{k+1}).$$

Noting that $\mu_{k+1} = \mu_k^\tau$, we have by (2.12)

$$\mu_{k+1}^\kappa \|g_{k+1}\|_0 \leq C_{s^*} \left\{ \mu_k^{\kappa\tau-(s^*-2-\beta)} C^k \|g_0\|_{s^*} + \theta_k \mu_k^{\kappa\tau+2} \|g_k\|_0 + \mu_k^{\kappa\tau+n+5} \|g_k\|_0^2 \right\}.$$

By (2.14) we get

$$(2.16) \quad \|g_{k+1}\|_0 \leq C_{s^*} \left\{ C^{k+1} \mu_k^{-2-n} \|g_0\|_{s^*} + \theta_k \mu_k^{2\kappa-2n-4} \|g_k\|_0 + \theta_k \mu_k^{2\kappa-n-1} \|g_k\|_0^2 \right\}.$$

Taking σ large enough, it follows that

$$\mu_{k+1}^\kappa \|g_{k+1}\|_0 \leq \frac{1}{4} \|g_0\|_{s^*} + d_k^2.$$

In the same way, using (2.13) and (2.14) we get

$$\mu_{k+1}^\kappa \theta_{k+1} \leq \frac{1}{4} \|g_0\|_{s^*} + d_k^2,$$

and then we have proved that

$$(2.17) \quad d_{k+1} \leq \frac{1}{4} \|g_0\|_{s^*} + d_k^2.$$

Now

$$(2.18) \quad g_0 = -G(w_0) = \frac{1}{\varepsilon} \check{f} \left(\varepsilon^2(x, y); \varepsilon^4 \sum_{i \leq n-1} \sigma_i(x_i^2 + y_i^2); 2\varepsilon^2 \sum_{i \leq n-1} \sigma_i(x_i, y_i) \right)$$

and $\check{f}(0; 0; 0) = 0$, so we can suppose that $\|g_0\|_{s^*} \leq 1$. Moreover, we assume that

$$\max(\mu_0^{2\kappa} \theta_0, \mu_0^{2\kappa} \|g_0\|) \leq \frac{1}{4 \max(1, K_{0,n+1})}.$$

Here $K_{0,n+1}$ is given by Lemma 1.3. By (2.15) and (2.17) we can conclude that $d_1 \leq \frac{1}{2} \|g_0\|_{s^*}$. By induction it is easy to see that we have

$$(2.19) \quad d_{k+1} \leq \frac{1}{2} \|g_0\|_{s^*},$$

which proves (2.10). \square

Proof of Theorem 1. We shall prove by induction that for some constant Γ

$$(2.20_k) \quad \|w_k\|_{2n+4} \leq \Gamma.$$

Since $w_0 = 0$, we may suppose that (2.20_k) is true for all $k = 0, 1, \dots, m$. By (2.4), and using the Gagliardo-Nirenberg inequality, we get

$$\|w_{m+1}\|_{s'} \leq \sum_{k=0}^m \|S_k \rho_k\|_{s'} \leq C \sum_{k=0}^m \|\rho_k\|_{s'} \leq C \sum_{k=0}^m \|\rho_k\|_{s^*}^{\frac{s'}{s^*}} \|\rho_k\|_0^{1-\frac{s'}{s^*}}.$$

By (1.10_{s*}), (2.8), and (2.9)

$$\|\rho_k\|_{s^*} \leq C (\|g_0\|_{s^*} + \|w_k\|_{s^*+4}) \leq C' \|g_0\|_{s^*} (1 + C^{k+1} \mu_k^\beta) \leq C_1^k \mu_k^\beta \|g_0\|_{s^*},$$

and by (2.10), $\|\rho_k\|_0 \leq C \|g_k\|_0 \leq C' \mu_k^{-\kappa} \|g_0\|_{s^*}$. It follows that

$$(2.21) \quad \|w_{m+1}\|_{s'} \leq C \sum_{k=0}^m C_1^k \mu_k^{\frac{s'}{s^*} \beta - \kappa (1 - \frac{s'}{s^*})} \|g_0\|_{s^*}.$$

We may fix the constants now. We first choose $\tau = \frac{4}{3}$, $\beta = 12 + \delta$, with $\delta > 0$ small, $\kappa = 12 + 3n + \delta$; then the first inequality in (2.14) is satisfied, and for the second one we have to take $s^* > 5n + 33$.

Let us go back to (2.21); then, taking $\frac{\varepsilon'}{s^*}$ sufficiently small and σ large, the series is convergent and we get $\|w_{m+1}\|_{s'} \leq C_{s^*} \|g_0\|_{s^*}$. The induction step is completed if we take $\varepsilon(s^*)$ so small that $C_{s^*} \|g_0\|_{s^*} \leq \Gamma$ (see (2.18)).

We have $w_m \rightarrow w$ in $H_s(\Omega)$ and

$$(2.22) \quad \|w\|_s \leq C_{s^*} \|g_0\|_{s^*}.$$

By (2.10), $g_m \rightarrow 0$; then, w is a solution of $G(w) = 0$ if $s \geq n+3$, and so $\phi = \sum_{i=0}^{n-1} \sigma_i |z_i|^2 + \varepsilon^5 w(\varepsilon^{-2}z)$ is a solution of the original Monge-Ampère equation,

which is plurisubharmonic if ε is small enough since $\varepsilon \left| \frac{\partial^2 w}{\partial z_i \partial z_j} \right| = O(\varepsilon)$. This completes the proof of Theorem 1. \square

3. EXISTENCE OF A C^∞ LOCAL SOLUTION

We shall use the result of C.J. Xu and C. Zuilly [12], [13], which we recall briefly. Let us consider a nonlinear partial differential equation

$$F(x, y, u, \nabla u, D^2 u) = 0,$$

where F is C^∞ . To any solution u we can associate the vector fields

$$X_j = \sum_k \frac{\partial F}{\partial u_{jk}} \partial_k.$$

Then we get

Theorem 3.1 ([12]). *Suppose $u \in C_{loc}^\rho(\Omega)$ with $\rho > \max(4, r+2)$ for some constant $r \geq 0$, and that the brackets of the X_j , up to the order r , span the tangent space at each point of Ω . Then u belongs to $C^\infty(\Omega)$.*

We can replace $\frac{\partial^2 \phi}{\partial z_i \partial \bar{z}_j}$ by the sum

$$\frac{1}{4} \left\{ \frac{\partial^2 \phi}{\partial x_i \partial x_j} + \frac{\partial^2 \phi}{\partial y_i \partial y_j} + i \frac{\partial^2 \phi}{\partial x_i \partial y_j} \right\};$$

so, by noting that $y_i = x_{i+n}$, $1 \leq i \leq n$, we obtain for our particular equation

$$\begin{aligned} X_i &= \phi^{ii} \frac{\partial}{\partial x_i} + \sum_{\substack{j=1 \\ j \neq i}}^n \frac{\phi^{ij} + \bar{\phi}^{i\bar{j}}}{2} \frac{\partial}{\partial x_j} + \sum_{\substack{j=1 \\ j \neq i}}^n \frac{i\phi^{ij} - i\bar{\phi}^{i\bar{j}}}{2} \frac{\partial}{\partial x_{j+n}}, \\ X_{i+n} &= \phi^{ii} \frac{\partial}{\partial x_{i+n}} + \sum_{\substack{j=1 \\ j \neq i}}^n \frac{\phi^{ij} + \bar{\phi}^{i\bar{j}}}{2} \frac{\partial}{\partial x_{j+n}} + \sum_{\substack{j=1 \\ j \neq i}}^n \frac{i\phi^{ij} - i\bar{\phi}^{i\bar{j}}}{2} \frac{\partial}{\partial x_j}. \end{aligned}$$

Hence

$$(3.1) \quad X_i = \varepsilon w_{n\bar{n}} A_i \partial_{x_i} + \varepsilon \sum_{l=1}^{2n} \sum_{j \neq i} w_{i\bar{j}} P_{ijl} (D^2 w) \partial_{x_l} \quad \text{for } i \notin \{n, 2n\},$$

and

$$(3.2) \quad \begin{cases} X_n = A_n \partial_{x_n} + \varepsilon \sum_{\substack{j \neq \delta \\ l}} w_{\delta \bar{j}} P_{\delta j l} (D^2 w) \partial_{x_l}, \\ X_{2n} = A_n \partial_{x_{2n}} + \varepsilon \sum_{\substack{j \neq \delta \\ l}} w_{\delta \bar{j}} P'_{\delta j l} (D^2 w) \partial_{x_l}. \end{cases}$$

Here $P_{\delta j l}$, $P'_{\delta j l}$ are polynomials in $D^2 w$, $A_n = \det (\delta_i^j \sigma_i + \varepsilon w_{i \bar{j}}; 1 \leq i, j \leq n-1)$ and A_l is the cofactor of $\sigma_l + \varepsilon w_{l \bar{l}}$ in A_n . In order to compute the Lie algebra generated by the X_i , we need some lemmas.

Lemma 3.2. *Let $\partial_{(x,y)}^\alpha \partial_u^l \partial_p^\beta f(0,0,0) = 0$ for all $|\alpha| + |\beta| + l \leq k-1$. Assume $w \in H^s$ with $s \geq n+k+1$. Then*

$$(3.3) \quad |\partial^\alpha w| \leq C_\alpha \varepsilon^{2k-1}, \quad \text{for all } |\alpha| \leq k.$$

Proof. Using (2.18), it follows from the hypothesis that the Taylor expansion of g_0 is on the form

$$(3.4) \quad \begin{aligned} g_0 &= \varepsilon^{4k-1} \sum_{|\alpha|+|\beta|+l=k} C_{l\alpha\beta} \varepsilon^{2l} (x,y)^\alpha (x',y')^\beta |x' + iy'|^{2l} \int_0^1 (1-\lambda)^k \\ &\times \partial_{(x,y)}^\alpha \partial_u^l \partial_p^\beta f \left(\lambda \varepsilon^2 (x,y); 2\lambda \varepsilon^4 \sum_{i \leq n} \sigma_i (x_i^2 + y_i^2); 2\lambda \varepsilon^2 \sigma_i (x',y') \right) d\lambda, \end{aligned}$$

and (3.3) follows from the inequality (2.22) and Lemma 1.3. \square

Lemma 3.3. *In addition to the assumption in Lemma 3.1, let*

$$\partial_{y_n}^k f(0,0,0) > 0.$$

Then, if ε is small enough,

$$(3.5) \quad \partial_{y_n}^k w_{n\bar{n}} \geq C \varepsilon^{2k-1}$$

with a constant C independent of ε .

Proof. Since $A_n = \Upsilon + \varepsilon \sum_{i,j} w_{i \bar{j}} A_{ij} (D^2 w)$, we rewrite the equation $\det (\phi_{i \bar{j}}) = f$ as

$$(3.6) \quad w_{n\bar{n}} + \varepsilon \sum_{i,j,k,m} w_{i \bar{j}} w_{k \bar{m}} A_{ijkm} (D^2 w) = \frac{f}{\varepsilon \Upsilon}.$$

Now

$$(3.7) \quad \begin{aligned} &\frac{1}{\varepsilon} f \left(\varepsilon^2 x; \varepsilon^4 \sigma_i x_i^2 + \varepsilon^5 w, (2\varepsilon^2 \sigma_i x_i + \varepsilon^3 \partial_{x_i} w)_{i \neq n, 2n}, \varepsilon^3 \partial_{x_n} w, \varepsilon^3 \partial_{x_{2n}} w \right) \\ &= g_0 + \varepsilon^4 w \Phi_1(x, \varepsilon, w, \nabla w) + \varepsilon^2 \sum_{p=2}^{2n+1} w_{p-1} \Phi_p(x, \varepsilon, w, \nabla w), \quad \Phi_j \in C^\infty. \end{aligned}$$

Applying to both sides of (3.6), (3.7) the operator $\partial_{y_n}^k$ and using (3.3), we get

$$\partial_{y_n}^k w_{n\bar{n}} \geq \frac{1}{\Upsilon} \partial_{y_n}^k g_0 - C_k (\varepsilon^{2k+3} + \varepsilon^{2k+1} + \varepsilon^{4k-1});$$

so,

$$(3.8) \quad \partial_{y_n}^k w_{n\bar{n}} \geq \frac{1}{\Upsilon} \varepsilon^{2k-1} \partial_{y_n}^k f(0,0,0) - C_k (\varepsilon^{2k+3} + \varepsilon^{2k+1} + \varepsilon^{4k-1}).$$

But $f(0,0,0) = 0$. Then $k \geq 1$, and (3.5) is proved. \square

Proof of Theorem 2. Without loss of generality we can assume that $\alpha^* = (0, \dots, 0, k)$ and $\partial_{y_n}^k f(0, 0, 0) > 0$. Choose s so big that $s \geq k + n + 4$. By means of Theorem 1 we can get a solution $w \in H^s$. Moreover, w satisfies (3.3) and (3.5).

* When $n > 2$ the polynomials B_{ijl} occurring in (3.1) and (3.2) are at least of degree 1. Thus, for each $i \in \{1, \dots, n-1, n+2, \dots, 2n-1\}$, using induction on the size of the bracket, one can prove that

$$(adX_{2n})^k(X_i) = \varepsilon \sigma_i \partial_{y_n}^k \omega_{n\bar{n}} A_n^k \partial_{x_i} + \varepsilon \sum_{i=1}^{2n} \sum_{\substack{|\alpha|+|\beta| \leq k \\ 1 \leq a, p, b, q \leq 2n}} C_{\alpha\beta abpq}^i \partial^\alpha \omega_{a\bar{b}} \bar{\partial}^\beta \omega_{p\bar{q}} \partial_{x_i},$$

where the $C_{\alpha\beta abpq}^i$ are still polynomials in $D^\gamma \omega$; $|\gamma| \leq k$.

From (3.3) and (3.5) the sum in the right hand side is $O(\varepsilon^{4k-1})$ while the coefficient of the first term is $\geq C_0 \varepsilon^{2k}$. It follows that the vector fields

$$\left[(adX_{2n})^k(X_i) \right]_{i=1, \dots, n-1, n+1, \dots, 2n-1}, X_n \text{ and } X_{2n}$$

span all the tangent space. Theorem 2 follows then from Theorem 3.1 and Lemma 1.3.

* If $n = 2$, a direct computation of the determinant of the vector fields

$$\left[(adX_4)^k(X_i) \right]_{i=1,3}, X_2 \text{ and } X_4$$

shows that it is equal to

$$\varepsilon^2 \sigma^{2k+2} (\partial_{y_n}^k \omega_{n\bar{n}})^2 + \varepsilon^3 \sum_{\substack{|\alpha|+|\beta|+|\gamma| \leq k \\ 1 \leq a, p, b, q, i, j \leq 2n}} C_{\alpha\beta abpq} \partial^\alpha \omega_{a\bar{b}} \bar{\partial}^\beta \omega_{p\bar{q}} \partial^\gamma \omega_{i\bar{j}},$$

which by (3.3) and (3.5) doesn't vanish, and we conclude in the same way.

* If $n = 1$, we don't need the hypothesis of Theorem 2 to conclude. \square

4. SUPPLEMENT

As noted by Bedford and Taylor in [2], if we set $|z_i| = e^{x_i}$ and $\psi(x_1, \dots, x_n) = \phi(e^{x_1}, \dots, e^{x_n})$, then it is easily checked that

$$\frac{\partial^2 \psi}{\partial x_i \partial x_j} = 4z_i \bar{z}_j \frac{\partial^2 \phi}{\partial z_i \partial \bar{z}_j}$$

and

$$\det \left(\frac{\partial^2 \psi}{\partial x_i \partial x_j} \right) = 4^n e^{2x_1} \dots e^{2x_n} \det \left(\frac{\partial^2 \phi}{\partial z_i \partial \bar{z}_j} \right);$$

so, we can state, as a consequence of Theorem 1.2,

Theorem 3 ([5], Theorem 1). *Let $f \in C^\infty$ be nonnegative near a point $Z^0 = (y_0, u_0, p_0) \in \mathbb{R}^n \times \mathbb{R} \times \mathbb{R}^n$ and $f(Z^0) = 0$. Then for any integer $s > n + 3$, the problem*

$$(4.1) \quad \det \left(\frac{\partial^2 \psi}{\partial y_i \partial y_j} \right) = f(y, u, \nabla u)$$

has a convex solution $u \in H^s$ in a neighborhood of y_0 , such that $u(y_0) = u_0$ and $\nabla u(y_0) = p_0$.

Theorem 4 ([5], Theorem 2). *Suppose in addition to the assumption in Theorem 3 that $\partial_y^\alpha \partial_u^l \partial_p^\beta f(y_0, u_0, \nabla u_0) = 0$ for all $|\alpha| + |\beta| + l \leq k - 1$ and there exists $\alpha^* \in \mathbb{N}^n$ such that $|\alpha^*| = k$ and $\partial_y^{\alpha^*} f(y_0, u_0, \nabla u_0) \neq 0$. Then (4.1) has a C^∞ convex local solution u in a neighborhood of y_0 .*

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